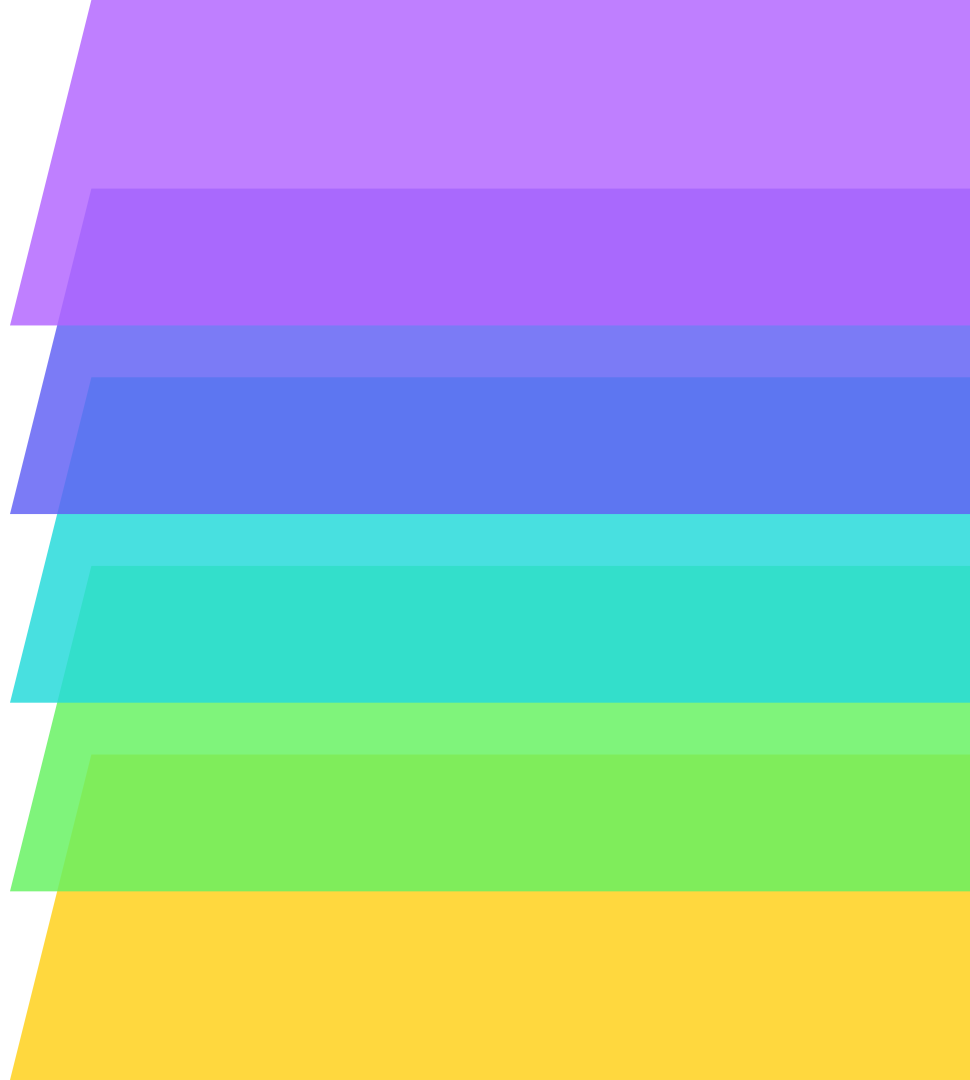


# Magnetic monopoles

Diego França de Oliveira



**1ST**

## Predicting Magnetic Monopoles

Quantization of charge, gauge symmetries and super massive monopoles

**2ND**

## Looking for magnetic monopoles

Detection experiments, astronomical bounds and the “monopole problem”

**3RD**

## Creating a magnetic monopole

MoEDAL experiment

**1ST**

# Predicting Magnetic Monopoles

# Maxwell's equation symmetry (in gaussian units)

In vacuum

$$\begin{aligned}\vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} &= 0 \\ \vec{\nabla} \cdot \vec{E} &= 0 \\ \vec{\nabla} \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} &= 0\end{aligned}$$

With charge

$$\begin{aligned}\vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} &= 0 \\ \vec{\nabla} \cdot \vec{E} &= 4\pi\rho \\ \vec{\nabla} \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} &= \frac{4\pi}{c} \vec{J}.\end{aligned}$$

# Maxwell's equation symmetry (in gaussian units)

In vacuum

$$\begin{aligned}\vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} &= 0 \\ \vec{\nabla} \cdot \vec{E} &= 0 \\ \vec{\nabla} \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} &= 0\end{aligned}$$

With charge

$$\begin{aligned}\vec{\nabla} \cdot \vec{B} &= 4\pi\rho_m \\ \vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} &= -\frac{4\pi}{c} \vec{J}_m \\ \vec{\nabla} \cdot \vec{E} &= 4\pi\rho \\ \vec{\nabla} \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} &= \frac{4\pi}{c} \vec{J}.\end{aligned}$$

# The absence of magnetic monopoles

Electrodynamics

$$\vec{B} = \vec{\nabla} \times \vec{A}$$

$$\nabla \cdot \vec{B} = \nabla \cdot (\vec{\nabla} \times \vec{A})$$

$$\nabla \cdot \vec{B} = 0$$

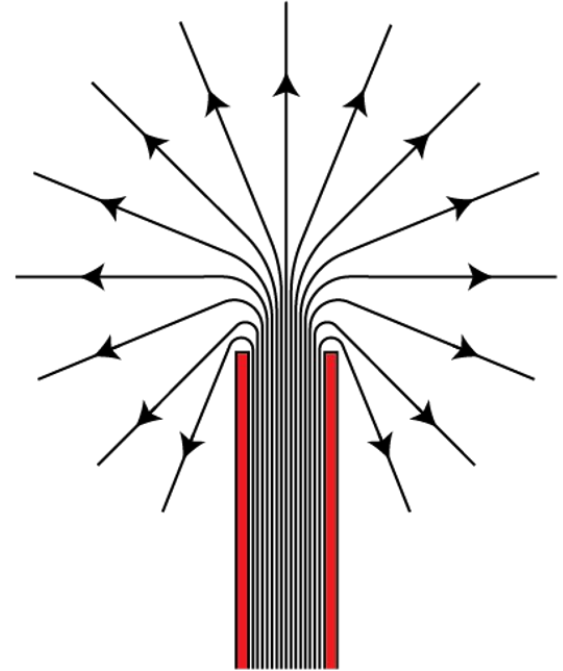
Quantum mechanics

$$|\psi_C\rangle = \exp\left(\frac{-iq}{\hbar c} \int \vec{A} \cdot d\vec{l}\right) |\psi_A\rangle$$

Potential fields forbid magnetic monopoles in ED and QM

# Dirac's argument (1931)

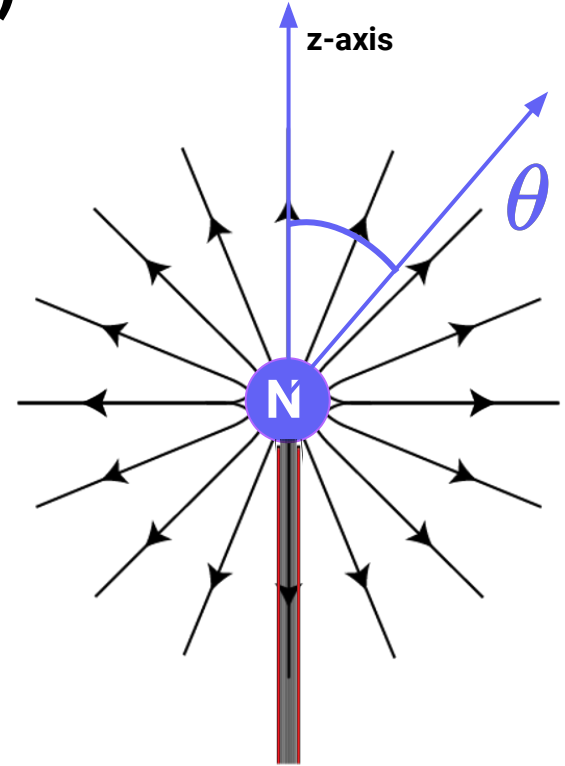
- A current through a solenoid produces magnetic field
- If the solenoid is infinitely long and thin, the field lines never close
- We have something similar to a magnetic monopole, but does the theory hold up?



# Dirac's argument (1931)

Magnetic field produced:  $\vec{B} = \frac{g}{r^2} \hat{r}$

A possible vector potential:  $\vec{A} = g \frac{1-\cos(\theta)}{r \sin(\theta)} \hat{\varphi}$



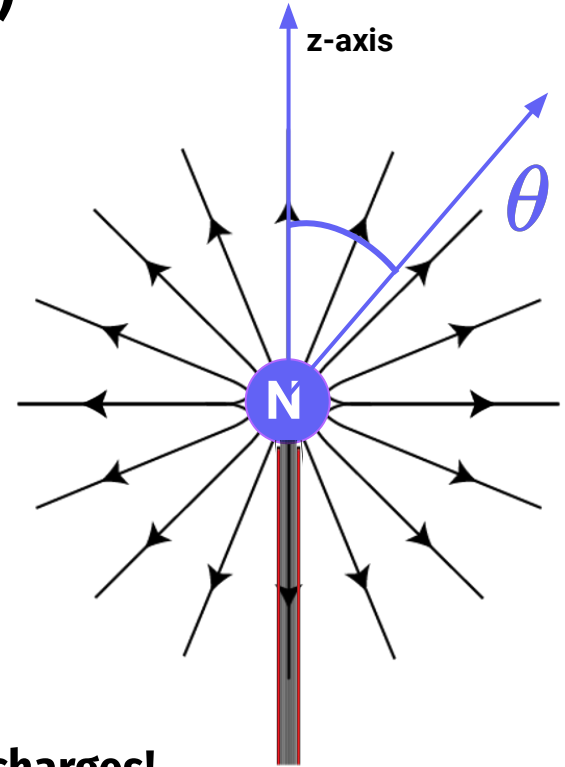


# Dirac's argument (1931)

**Flux of magnetic field:**

$$\begin{aligned}\Phi_B &= \int_A \vec{B} \cdot d\vec{S} = \int_\Gamma \vec{A} \cdot d\vec{l} \\ &= 2\pi g(1 - \cos(\theta))\end{aligned}$$

$$\Phi_B(\theta = 2\pi) = 4\pi g$$



**Similar to the gauss law for electric charges!**

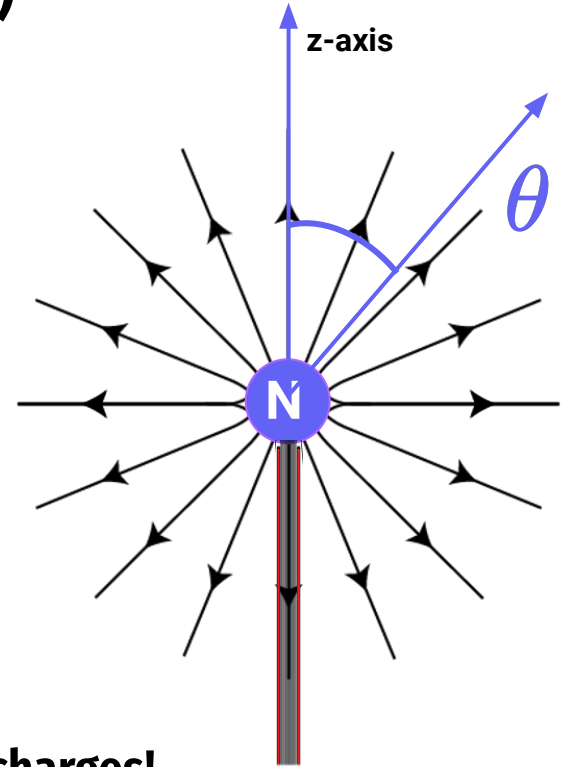
$$\Phi_E(\theta = 2\pi) = 4\pi q$$

# Dirac's argument (1931)

Flux of magnetic field:

$$\begin{aligned}\Phi_B &= \int_A \vec{B} \cdot d\vec{S} = \int_\Gamma \vec{A} \cdot d\vec{l} \\ &= 2\pi g(1 - \cos(\theta))\end{aligned}$$

$$\Phi_B(\theta = 2\pi) = 4\pi g$$



Similar to the gauss law for electric charges!  
**But is mathematically wrong**

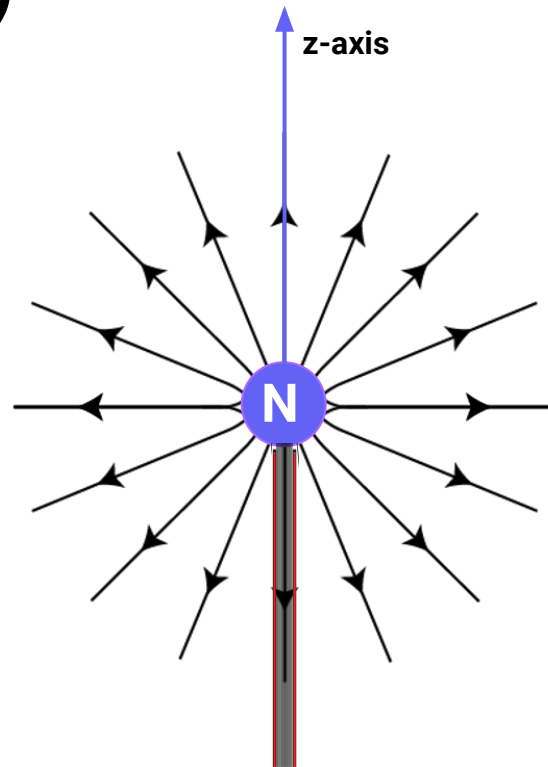
# Dirac's argument (1931)

Flux of magnetic field:

$$\Phi_B = \begin{cases} 2\pi g(1 - \cos(\theta)); & \theta < \pi \\ 0 & ; \theta = \pi \end{cases}$$

$$\vec{B} = \nabla \times \vec{A} - 4\pi g \Theta(-z) \delta(x) \delta(y) \hat{z}$$

**Dirac string singularity**  
Physical or mathematical?



# Dirac's argument (1931)

Is the dirac string detectable?

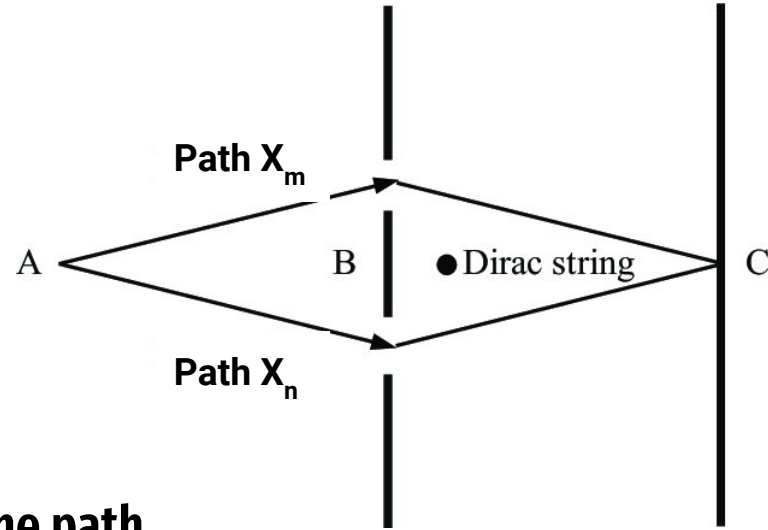
By the Aharonov-bohm effect:

$$|\psi_C\rangle = \exp\left(\frac{-iq}{\hbar c} \int \vec{A} \cdot d\vec{l}\right) |\psi_A\rangle$$

The phase change will be

$$\theta = \frac{q}{\hbar c} \int \vec{A} \cdot d\vec{l}$$

Measuring the phase in C allows one to know the path.



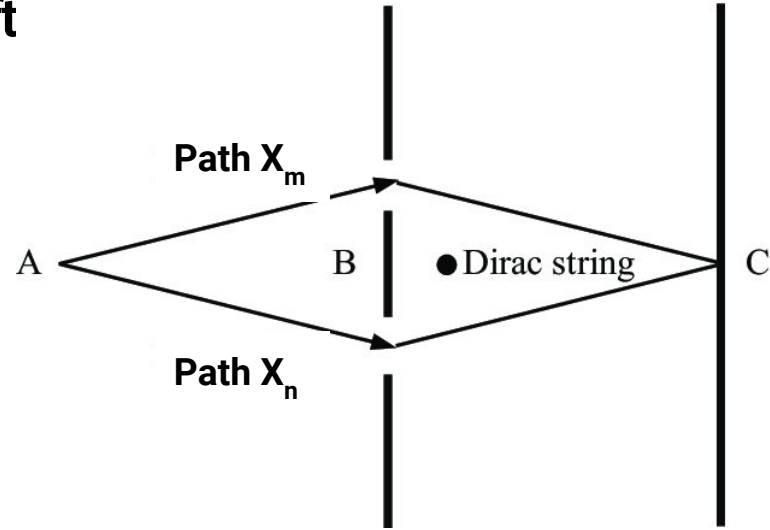
# Dirac's argument (1931)

It's possible to calculate the difference in phase shift  
Between paths.

$$\Delta\theta = \frac{-q}{\hbar c} \int_{A \rightarrow X_m \rightarrow C} \vec{A} \cdot d\vec{l} - \frac{-q}{\hbar c} \int_{A \rightarrow X_n \rightarrow C} \vec{A} \cdot d\vec{l}$$

$$\Delta\theta = \frac{-q}{\hbar c} \int_A \vec{B} \cdot d\vec{S}$$

$$\Delta\theta = \frac{-q}{\hbar c} 4\pi g$$



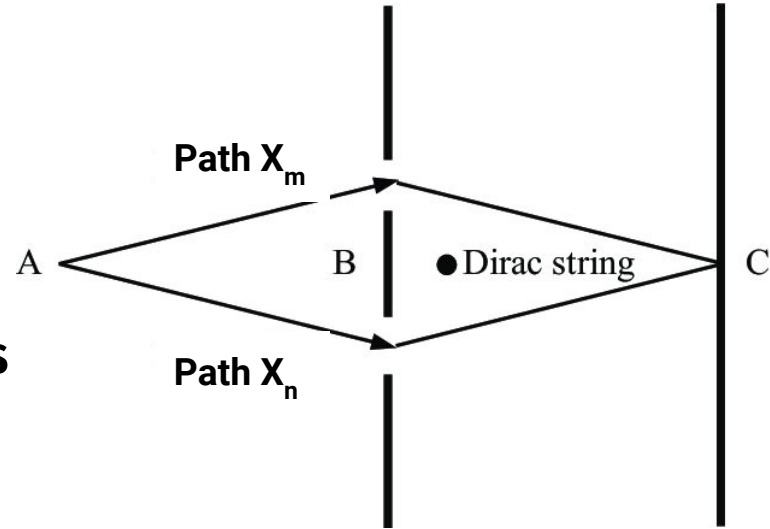
# Dirac's argument (1931)

There is a condition where is undetectable:

$$\frac{-q}{\hbar c} 4\pi g = n2\pi$$

$$\frac{-qg}{\hbar c} \in \mathbf{Z}$$

If **charge is quantized**, then magnetic monopoles are allowed!

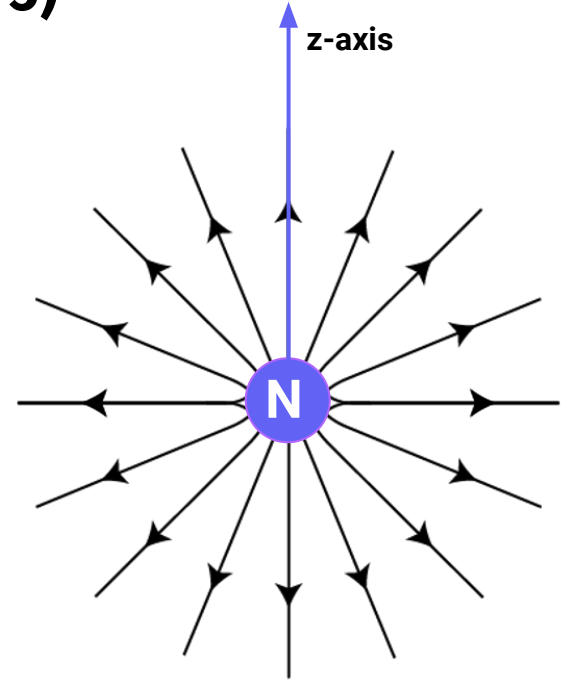


# Removing the string (1975)

Charge quantization comes from the dirac string situation. Can we remove the string while getting the quantization condition?

$$\nabla \cdot \vec{B} = \nabla \cdot (\vec{\nabla} \times \vec{A})$$

Not using a vector potential...Let's try to use more



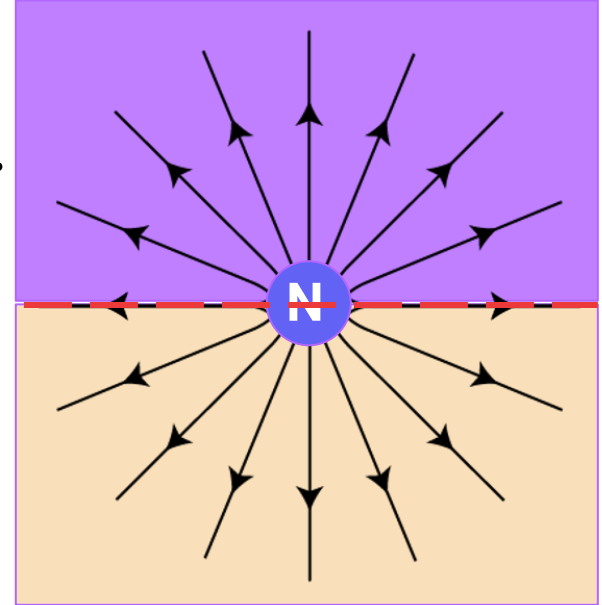
# Removing the string (1975)

Let's assume two vector potential in two different regions.

$$\vec{A}_{sup} = g \frac{1 - \cos(\theta)}{r \sin(\theta)} \hat{\varphi}; 0 \leq \theta \leq \frac{\pi}{2}$$

$$\vec{A}_{inf} = -g \frac{1 + \cos(\theta)}{r \sin(\theta)} \hat{\varphi}; \frac{\pi}{2} \leq \theta \leq \pi$$

$$\left. \begin{array}{l} \vec{A}_{sup} \\ \vec{A}_{inf} \end{array} \right\} \vec{B} = \frac{g}{r^2} \hat{r}$$





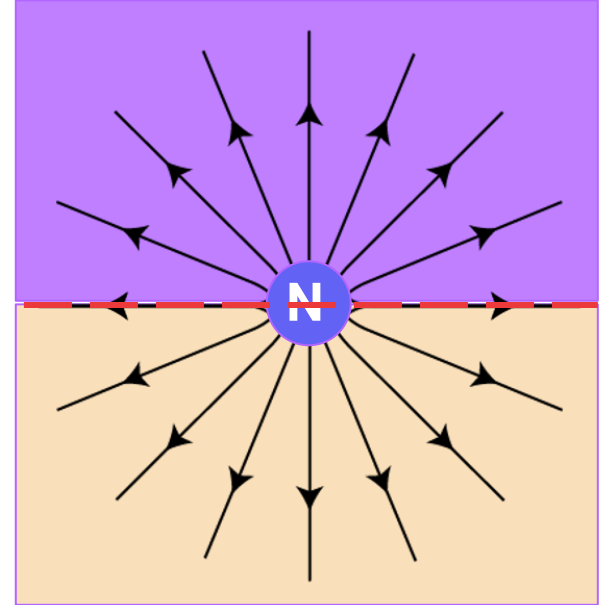
# Removing the string (1975)

In the equator, is mandatory that:

$$\vec{\nabla}\chi = \vec{A}_{sup} - \vec{A}_{inf}$$

$$\vec{\nabla}\chi = -\frac{2g}{r\sin(\theta)}\hat{\varphi}$$

$$\chi = 2g\varphi$$



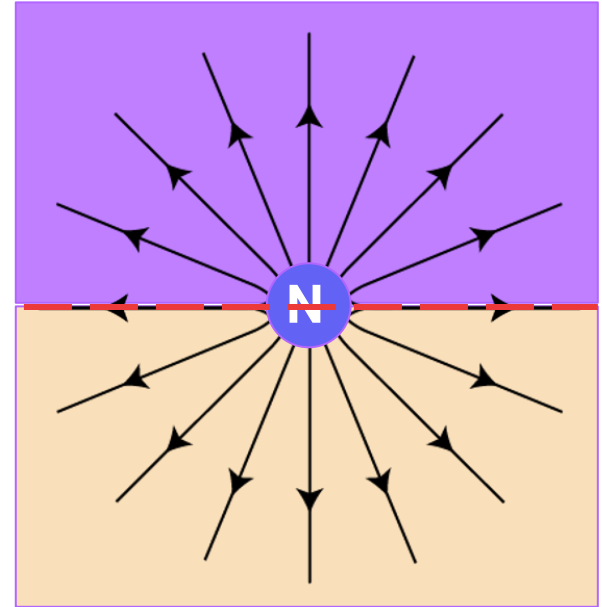
# Removing the string (1975)

The same gauge field must be valid in QM:

$$U_g = \exp\left(\frac{-iq\chi(\varphi=0)}{\hbar c}\right) = \exp\left(\frac{-iq\chi(\varphi=2\pi)}{\hbar c}\right)$$

$$U_g = \exp(0) = \exp\left(\frac{-i4qg\pi}{\hbar c}\right)$$

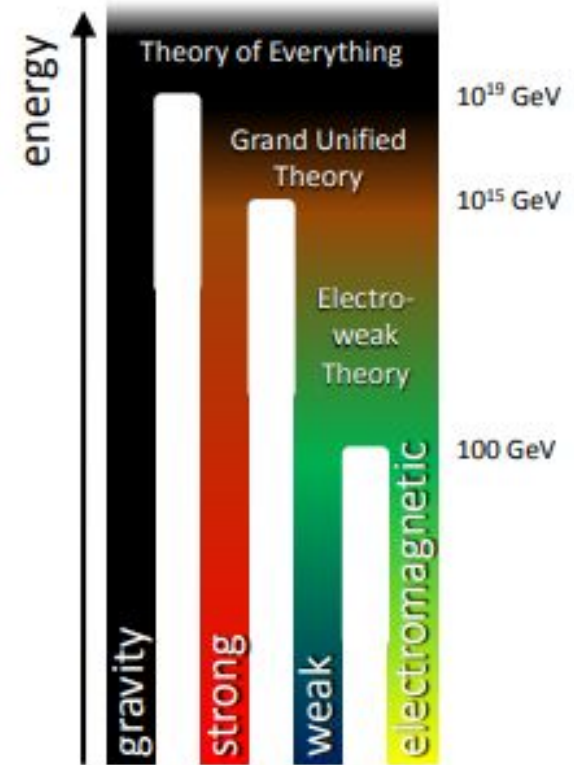
$\frac{2qg}{\hbar c} \in \mathbf{Z}$  The charge quantization is a consequence of gauge symmetries!



# 'T Hooft-Polyakov Monopole(1975)

In simple terms

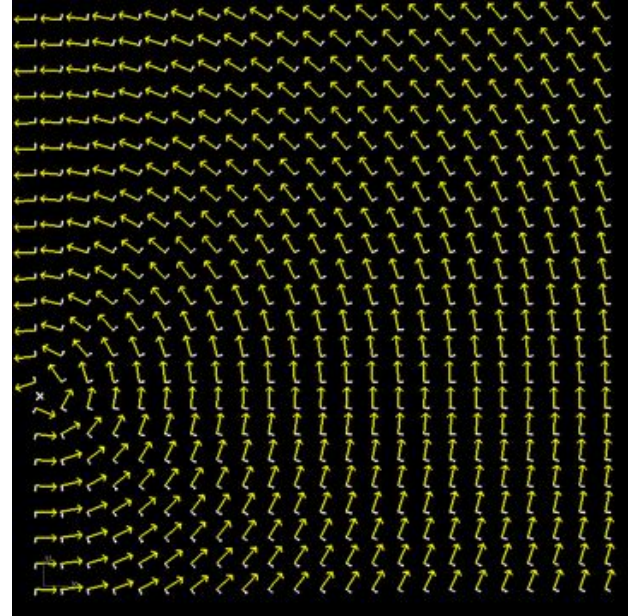
- The gauge field theory allows one to study the fundamental forces in a unified way in high energies.
- What separates the two symmetries, at lower energies, is the higgs field.



# 'T Hooft-Polyakov Monopole(1975)

In simple terms

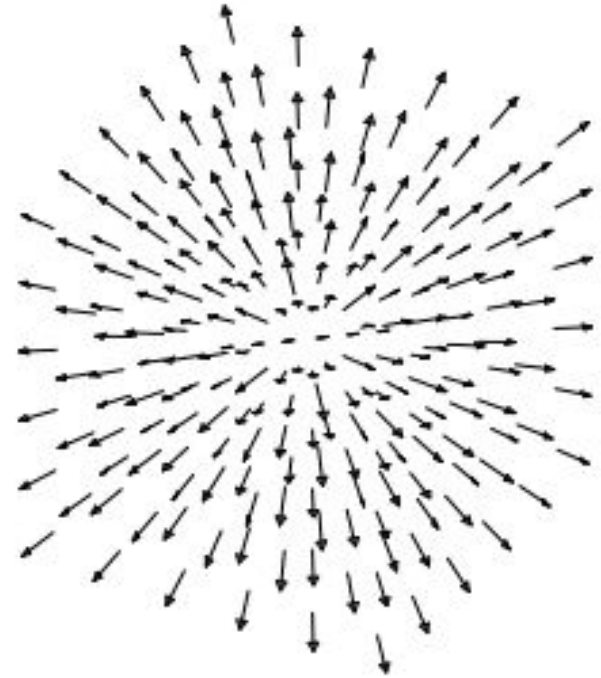
- The higgs field in Grand Unified Theories can be *represented as a vector field* with two characteristics.
  - Continuous and invariant under smooth changes (similar to vector field gauge invariance)
  - Non-zero everywhere in vacuum



# 'T Hooft-Polyakov Monopole(1975)

In simple terms

- The “Hedgehog” configuration is a interesting situation.
- Field is zero in the center, which can't be removed by smooth changes.
  - It cannot be vacuum: there is a massive particle in the center ( $E \sim 10^{15} \text{ GeV}$ )
- In this situation, the magnetic field can be determined as
$$\vec{B} = \frac{g}{r^2} \hat{r}$$
- Gauge theories inevitably predict magnetic monopoles!



**2ST**

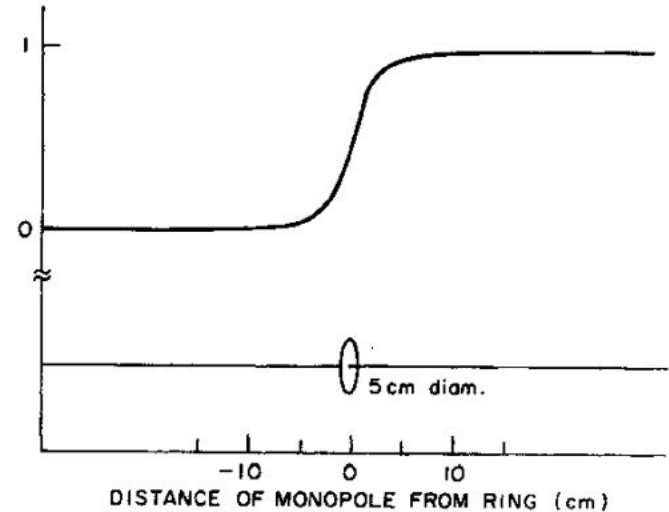
**Looking for magnetic monopoles**

# Blas Cabrera's experiment

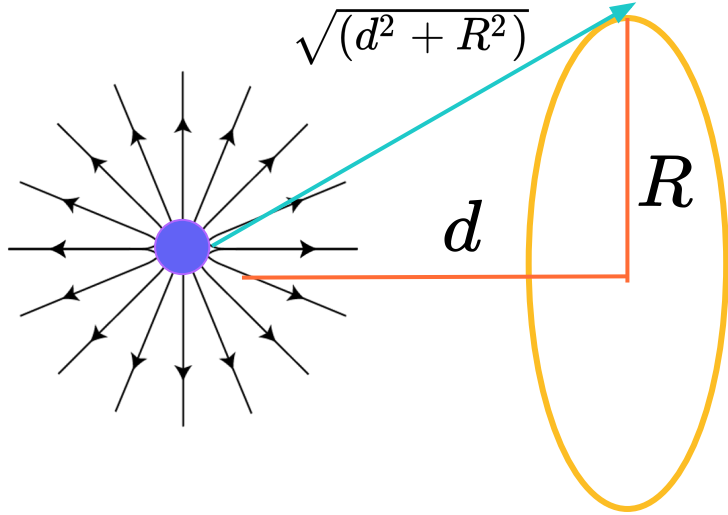
- Monopoles were never seen on earth, but is possible to look for them in cosmic radiation
- If a monopole passes through a superconductive coil, it will produce a current

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} - \frac{4\pi}{c} \vec{j}_m$$

$$\oint \vec{E} \cdot d\vec{l} = -\frac{1}{c} \frac{\partial \Phi_m}{\partial t} - \frac{4\pi}{c} \vec{I}_m$$



# Blas Cabrera's experiment



$$\Phi_B = \int_A \vec{B} \cdot d\vec{S}$$

$$\Phi_B = 2\pi g \left[ 1 - \frac{d}{\sqrt{(d^2 + R^2)}} \right]$$

- Taking  $t=0$  as the time  $d=0$ :

- $t < 0$

$$\Phi_B = 2\pi g \left[ 1 + \frac{vt}{\sqrt{((vt)^2 + R^2)}} \right]$$

- $t > 0$

$$\Phi_B = 2\pi g \left[ -1 + \frac{vt}{\sqrt{((vt)^2 + R^2)}} \right]$$

$$\Phi_B = 2\pi g \left[ 1 - 2\Theta(t) + \frac{vt}{\sqrt{((vt)^2 + R^2)}} \right]$$



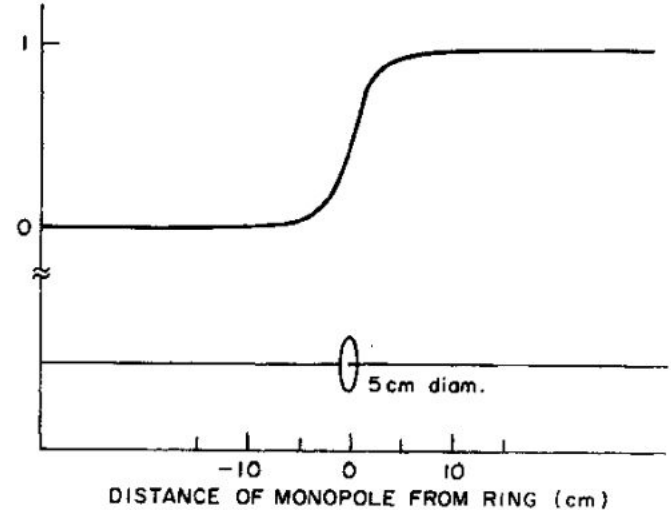
# Blas Cabrera's experiment

$$-LI_e(t) = -\frac{1}{c}\Phi_m - \frac{4\pi}{c}g\Theta(t) \quad \Phi_B = 2\pi g\left[1 - 2\Theta(t) + \frac{vt}{\sqrt{(vt)^2 + R^2}}\right]$$

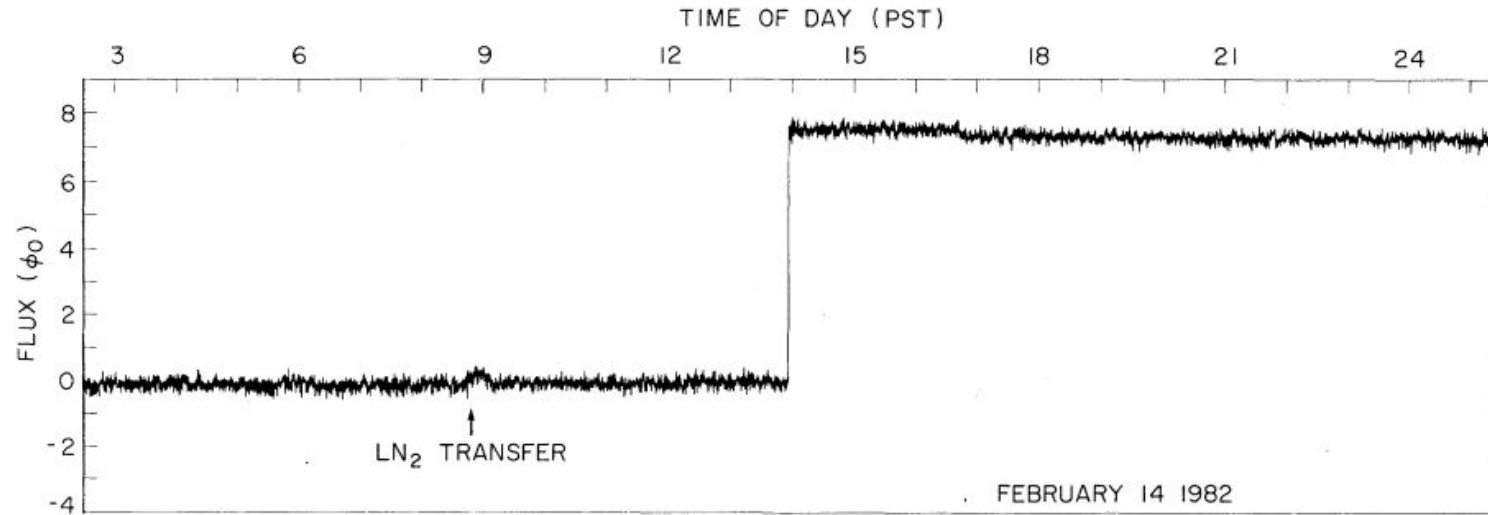
$$I(t) = \frac{2\pi g}{Lc} \left( 1 + \frac{vt}{\sqrt{(vt)^2 + R^2}} \right)$$

$$\lim_{t \rightarrow +\infty} I(t) = \frac{4\pi g}{L}$$

$$\lim_{t \rightarrow -\infty} I(t) = 0$$



# Blas Cabrera's experiment



# Other experiments



**Monopole, Astrophysics and Cosmic Ray Observatory (MACRO)**

- **MACRO operated from 1989 to 2002 with a detection area of  $10000\text{m}^2$ , never detecting one.**
- **No one ever saw any evidence of monopoles in space after 1982!**
- **If monopoles exist, would we be able to find them?**

# The poynting theorem with magnetic current

**Poynting vector:**

$$\vec{S} = \frac{c}{4\pi} \vec{E} \times \vec{B}$$

$$\nabla \cdot \vec{S} = -\frac{1}{4\pi} \left[ \vec{B} \cdot \frac{\partial \vec{B}}{\partial t} + \vec{E} \cdot \frac{\partial \vec{E}}{\partial t} \right] - \vec{B} \cdot \vec{j}_m - \vec{E} \cdot \vec{j}_e$$

**Field's energy density:**

$$u = \frac{E^2 + B^2}{8\pi}$$

$$\frac{\partial u}{\partial t} = \frac{1}{4\pi} \left[ E \frac{\partial E}{\partial t} + B \frac{\partial B}{\partial t} \right]$$

$$\nabla \cdot \vec{S} + \frac{\partial u}{\partial t} = -\vec{B} \cdot \vec{j}_m - \vec{E} \cdot \vec{j}_e$$

**The higher the currents, the  
faster the energy drops**

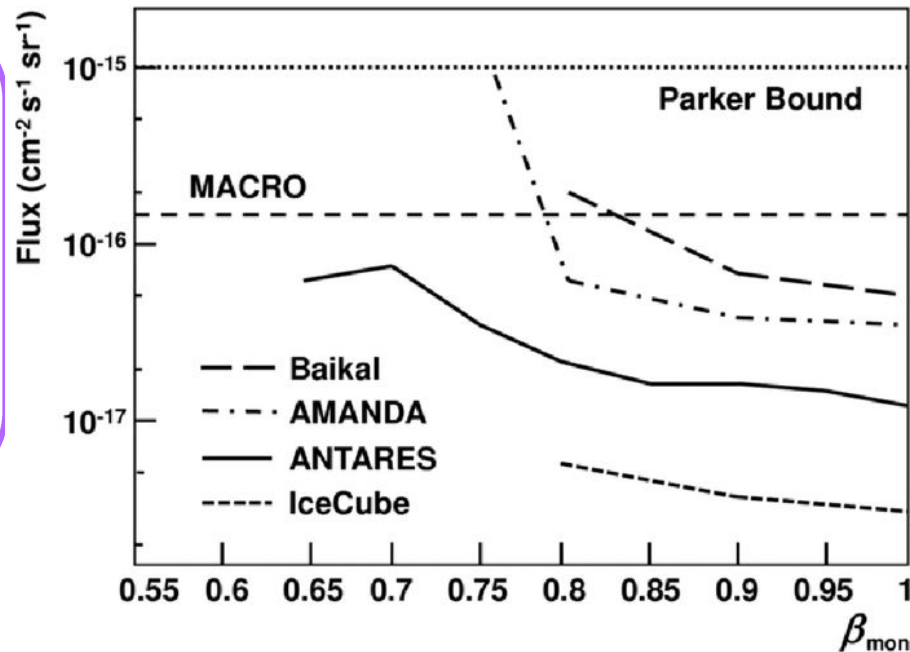
# Astrophysical bounds

$$\nabla \cdot \vec{S} + \frac{\partial u}{\partial t} = -\vec{B} \cdot \vec{j}_m - \vec{E} \cdot \vec{j}_e$$

As the galaxy has a longing magnetic field of  $3\mu G$  this drop cannot be so fast around here, so there is a limit to the magnetic current.

The flux  $F$  of magnetic field can be estimated. By Parker's bound, it would be

$$F \leq 10^{-15} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$



# Astrophysical bounds

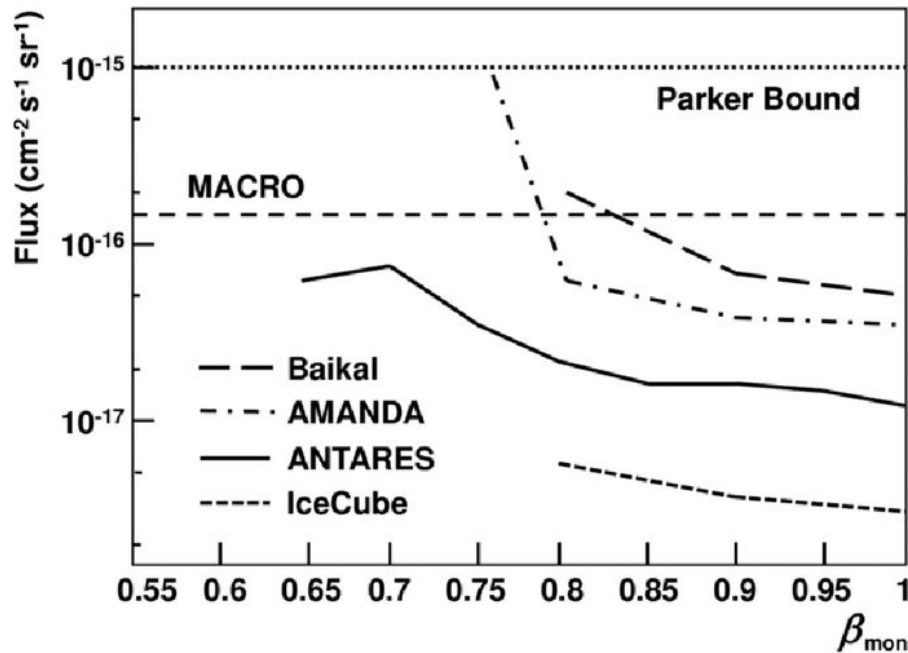
$$F \leq 10^{-15} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$

Considering a detector the size of MACRO (10000m<sup>2</sup>), the rate of passing monopoles would be.

$$\frac{N}{t} \leq 1.5 \cdot 10^{-18} \text{ s}^{-1}$$

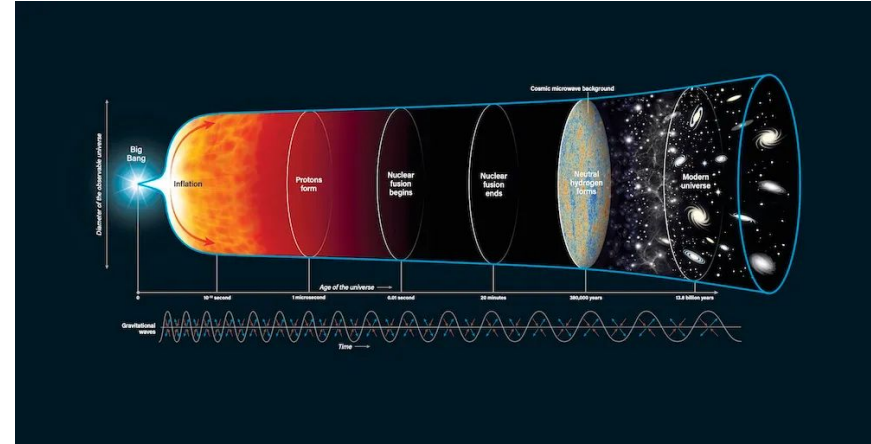
$$\frac{N}{t} \leq 5 \cdot 10^{-3} \text{ billion years}^{-1}$$

One every 20 billion years!?



# Astrophysical bounds

- If they were as massive as expected ( $E \sim 10^{15} \text{ GeV}$ ), they would have formed in the Big Bang.
- To compensate for their gravitational attraction, we need **cosmic inflation!** (“The **monopole problem**”)
- If that’s true, we will never find a monopole.
- Can we make one?



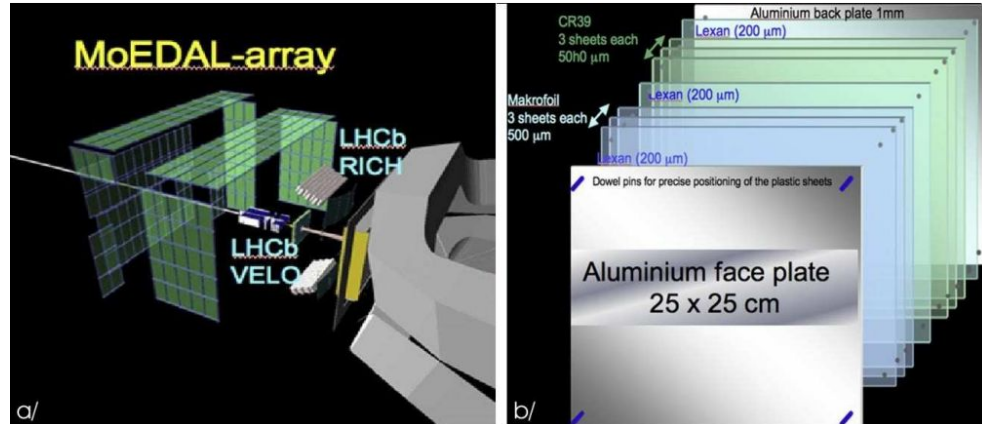
**3ST**

**Creating a magnetic monopole**



# MoEDAL

- Although not predicted, Intermediate Mass Monopoles are still allowed.
- The Monopole and Exotics Detector at the LHC(MoEDAL) , with collision energies of 8TeV, has been trying to create a monopole since 2009.



**4ST**

**Conclusion**

# Conclusion

- **Magnetic monopoles are allowed under ED and QM**
- **Gauge field theories expect them at higher energies**
- **We never find one, and probably never will...**
- **MoEDAL is the most modern experiment currently trying to prove their existence**

## Nobel prizes



**Paul Dirac(1933)**  
**Antimatter prediction**



**Carl Anderson(1936)**  
**Discovery of the positron**

# References

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# Thank you!

Special thanks: bruno Trebbi